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Fabrication and Spectroscopic Properties of Heavily Pr³⁺ Doped Selenide Chalcogenide Glass and Fiber for Mid-infrared Fiber Laser

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Abstract: In order to develop a high gain medium for fiber lasers operating at 3–5 μm waveband, 0–0.4% (in weight) Pr³⁺ ions doped Ge₁₂As_{20.8}Ga₄Se_{63.2} selenide chalcogenide glasses were prepared and the 0.2% (in weight) Pr³⁺ ions doped one was successfully drawn into step-index double-cladding fiber with the lowest loss of 2.95 dB/m@6.58 μm by a multistage rod-in-tube method. The electron-probe measure microanalysis (EPMA), X-ray diffraction (XRD), differential scanning calorimeter (DSC), field emission transmission electron microscope (FE-TEM), transmission and mid-infrared fluorescence spectra were carried out to analyze the dispersion of Pr³⁺ ions in glass, the impurity contents, thermal and optical changes caused by the Pr³⁺ ions' introduction. By analyzing the absorption and emission measurements of the serial glasses with the Judd-Ofelt theory, the Judd-Ofelt strength parameters, transition probabilities, excited state lifetime, branching ratios, and emission cross-sections were also calculated. This selenide chalcogenide glass has high Pr³⁺ ions' solubility and emission characteristic, good thermal stability and fiber forming performance, indicating that it has potential to be used as mid-infrared laser working medium.

Key words: chalcogenide glass; rare-earth ions; spectroscopy; mid-infrared fluorescence; infrared fiber

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用于中红外光纤激光器的高 Pr³⁺ 掺杂硒化物硫系玻璃和 光纤制备及其光谱特性

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摘要: 为了研发用于 3~5 μm 波段光纤激光器的增益介质, 制备了重量百分比为 0~0.4% 不同浓度 Pr³⁺ 离子掺杂的 Ge₁₂As_{20.8}Ga₄Se_{63.2} 硒化物硫系玻璃。通过多级棒管法, 重量百分比为 0.2% 的 Pr³⁺ 离子掺杂玻璃被成功拉制成阶跃型双包层光纤, 损耗最低为 2.95 dB/m (位于 6.58 μm 处)。采用电子探针显微分析 (EPMA)、X 射线衍射 (XRD)、差示扫描量热 (DSC)、场发射透射电子显微镜 (FE-TEM)、透射光谱和中红外荧光光谱分析了玻璃中

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Pr³⁺离子的分散性、杂质含量以及Pr³⁺离子引入引起的热、光学性质变化。通过玻璃的吸收和发射光谱并结合Judd-Ofelt理论,计算了Judd-Ofelt强度参数、辐射跃迁几率、荧光寿命、荧光分支比和受激发射截面。这种硒化硫系玻璃具有较高的Pr³⁺离子溶解度和中红外发光特性、良好的热稳定性和成纤性能,表明其具有作为中红外激光工作介质的潜力。

关 键 词: 硫系玻璃; 稀土离子; 光谱学; 中红外荧光; 红外光纤

1 Introduction

Rare-earth (RE) ions doped fiber lasers operating at mid-infrared (MIR) regions are of great interests due to their potential applications in fields of chemical sensing, air pollution monitoring, medical surgery, national defense, and *etc*^[1-4]. Chalcogenide glasses (ChGs), as hosts for RE ions, have their unique advantages. The ChGs' phonon energies are as low as 250–425 cm⁻¹, especially in which the selenide glasses have lower phonon energies of 250–300 cm⁻¹[5]. This results in lower probabilities of multi-phonon transitions for RE ions, making it easier to achieve higher efficient MIR transitions of RE ions.

It is known that the RE ions doped fluoride fibers have achieved high-power laser output in 3 μm waveband^[6-7], and compared with fluoride ones, ChG fibers have no obvious advantage for laser output at below 4 μm. Because of this, in the past two decades, people paid much attention to realize laser output at ≥4 μm wavelength in ChG fibers^[8-11]. However, there has been little experimental progress in the research of ChGs' fibers in laser output. There are three main reasons including the low RE ions' solubility, the still high multi-phonon relaxation rate of RE ions and the severe fiber loss originating from impurities' absorption^[12-13].

Recently, a longest laser output wavelength as well as 5.14 μm was reported in a 0.05% (in weight, the same below) Ce³⁺-doped Ge₁₅As₂₁Ga₁Se₆₃ fiber with a 9 μm-core-diameter pumped by a 4.15 μm quantum cascade laser at room temperature^[14]. Ge-As-Ga-Se (GAGSE) glass system is interested because of its two advantages. Firstly, compared with the stubborn S—H bonds in sulfide glasses, Se—H bonds are easier to be eliminated because of lower polarity and bond energy (339 kJ/mol for S—H

and 314 kJ/mol for Se—H). Secondly, the conventional As—Se system glasses have strong covalence, resulting in the tendency of clustering of RE ions and quenching of fluorescence. Based on the original glass network, the introduction of Ge can improve the compactness of glass grid and increase the laser damage threshold^[15-16], which means the glass can be pumped by a higher power laser. However, there are a large number of Se—Se bridging bonds and Ge—Se covalent bonds in Ge—As—Se, which cannot compensate the high coordination number of RE ions, resulting in the low RE's solubility. Ga is further introduced to compensate for negative local environment of positive RE ions by forming Ga—Se—RE³⁺^[9]. Introduction of Ga is effective to weaken the aggregation of RE ions in glass matrix.

Several researchers have conducted many studies of RE ions doped GAGSE glasses and fibers. In 2014, Tang *et al.* prepared a 0.05% Pr³⁺ ions doped GAGSE fiber by the extrusion method, and the lowest loss of fiber is 2.8 dB/m at 6.65 μm^[8]. In 2018, Liu *et al.* prepared the Pr³⁺ ions doped Ge₁₀As₂₄Ga₄Se₆₂ glass and fiber with an optimal doping concentration of 0.2%. The 250/125 μm (cladding diameter/core diameter), 0.2% doped step-index fiber exhibited the lowest loss of 3.5 dB/m at 6.9 μm. The fiber loss is about 10 dB/m at 2.9 μm caused by O—H bonds and 7 dB/m at 6.3 μm caused by H₂O molecule^[9]. In order to remove impurities and achieve low-loss GAGSE fibers, Shiryaev *et al.* reported a special chemical reaction device to transfer the GaI₃ vapor stream into the reactor for preparation of Ge₁₆As₁₇Ga₃Se₆₄ glasses^[17]. This provides a practical idea for the preparation and purification of Ga-containing GAGSE glasses. In 2018, they successfully fabricated a low-loss Ge₂₅As₁₅Ga₃Se₅₇ single-index fiber by using metallic aluminum and TeCl₄ for purification.

The minimum fiber loss is 0.6 dB/m at 6.7 μm and the optical loss is about 5 dB/m at 4.58 μm caused by Se—H bonds^[18]. However, there are still few reports about fiber laser output of Pr³⁺ ions. This is mainly because the severe absorption of impurities and weak fluorescence intensity in Pr³⁺ doped ChG fibers. So far, the doping concentration of Pr³⁺ ions in ChG fibers reported is relatively low (usually 0.5%). Therefore, the successful fabrication of a heavily Pr³⁺ doped CHG fiber with low impurity absorption has become an urgent requirement.

In this work, a modified glass composition of Ge₁₂As_{20.8}Ga₄Se_{63.2} was designed and a series of glasses doped with different concentration of Pr³⁺ ions were prepared. This composition glass has higher Pr³⁺ ions solubility (at least 0.4% without crystallization), and the 0.2% Pr³⁺-doped glass was drawn into fiber with double-cladding structure. For reducing content of impurities (oxygen, hydrogen, carbon, silicon, and *etc.*), the raw materials were distilled and the glasses were introduced with chemical impurity removal reagent. The absorption and emission spectra of bulk glasses were investigated and Judd-Ofelt intensity parameters were calculated. Broadband MIR luminescence of Pr³⁺ ions in the region of 3.3–5.5 μm was observed, and the stimulated emission cross-section, radiation lifetime and other parameters were calculated.

2 Experiment

2.1 Glass Preparation and Fiber Drawing

Ge₁₂As_{20.8}Ga₄Se_{63.2} glass samples doped with 0, 0.05%, 0.1%, 0.2%, 0.3% and 0.4% in weight of Pr³⁺ ions (denoted as GAGSE-undoped, GAGSE-0.05, GAGSE-0.1, GAGSE-0.2, GAGSE-0.3 and GAGSE-0.4) were prepared by melt quenching method. The raw materials of Ge (6N, grains, Nanjing Germanium Co., Ltd., RRC), Ga (7N, grains, Aladdin Industrial Co. Ltd., RRC) and Praseodymium powder (3N, powder, Aladdin Industrial Co. Ltd., RRC) were respectively packed in ampoules and heated to remove volatile impurities. The As (7N, grains, Mount Emei Jiamei High Pure Material

Co. Ltd., R. R. C.) and Se (6N, Aladdin Industrial Co. Ltd., RRC) were purified by a multi-stage distillation process. All of the raw materials were weighed by an analytical balance (Sartorius, resolution of 0.001 g) in the glovebox (N₂-filled, O₂ $\leq 10^{-8}$, H₂O $\leq 10^{-8}$), then loaded Ge, As, Se into a pre-treated quartz ampoule (800 °C for 12 h, under vacuum of 10⁻³ Pa). The ampoule was sealed under vacuum (10⁻³ Pa), and initially melted at 900 °C for 10 h to homogenize host glasses. The initial melted sample was transferred to a customized distillation ampoule, and Al foil (5N, Aladdin Industrial Co. Ltd., RRC) was introduced as deoxidizer. The distilled compound was then mixed with Ga and Pr powder and melted at 850 °C for 8 h. Subsequently, the quartz ampoule containing the molten glass was quenched, annealed near the glass transition temperature, and slowly cooled to the room temperature. The samples were cut into disks with a diameter of 12 mm and thickness of 3 mm, and mirror polished on both sides for following measurements.

The glass sample which has the best optical and thermal properties was selected to be drawn into fiber by a multistage rod-in-tube method (see Fig. 1). In order to achieve single-mode transmission and control the diameter of fiber easily during the fiber drawing, a double-cladding structure was designed. The cladding tubes were prepared by the rotational casting method. After the core rod and inner-cladding tube were combined, they were pre-drawn into a slim rod (external diameter $d=2.6$ mm). Then we assembled this preform with a thick outer-cladding tube which has a small hole through the center (internal diameter $d'=3.0$ mm), then coated it with FEP tube (external diameter $d=15.0$ mm). Finally, the single-mode, double-cladding fiber with a diameter of 125 μm was drawn from the assembly. The optical fiber draw tower (SG Control, UK; MT-600N) is equipped with a built-in furnace which heated the preform to 440 °C at a heating rate of 2 °C/min and kept this temperature to draw the optical fiber. During the drawing process, the He was introduced for atmosphere protection.

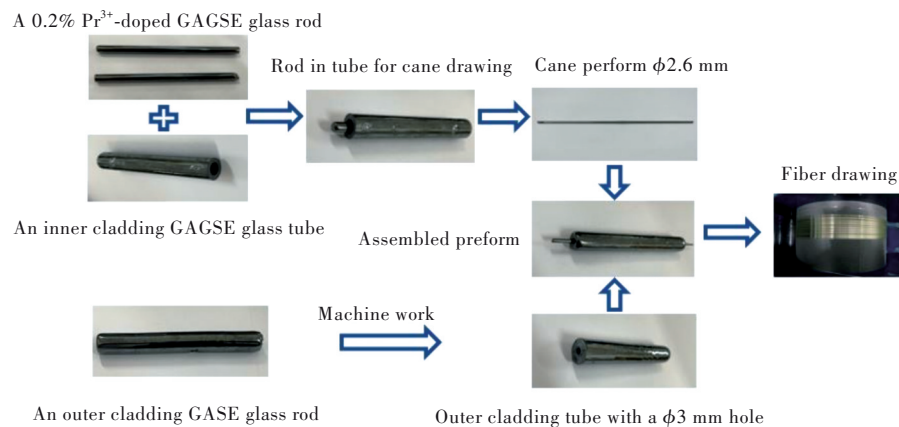


Fig. 1 The fabrication process of 0.2% Pr^{3+} -doped, single-mode and double-cladding $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ glass fiber.

2.2 Physical and Optical Measurements

To figure out the composition deviation of the glass samples caused by the chemical purification treatment, the electron-probe measure microanalysis (EPMA, JEOLJXA-8230) was carried out. All samples were analyzed for the presence of crystalline phases by X-ray diffraction (XRD) (Bruker D8 ADVANCE with Cu radiation, a power of 40 kV, 40 mA, a step size of $0.02^\circ (2\theta)$ and step time of 0.1 s). The thermal characteristics of the samples were analyzed by a Differential Scanning Calorimeter (NETZSCH DSC 404 F3). 15 mg of glass powder was put into an aluminum crucible and heated to 400°C under the protection of flowing atmosphere (N_2) with a heating rate of $10^\circ\text{C}/\text{min}$. The Field Emission Transmission Electron Microscope (FE-TEM) (FEI-Talos F200S/FEI Tecnai G2 F20 X-Twin) was used to verify the amorphous properties of the glass and the dispersion of RE ions with an Energy Dispersive Spectrometer (EDS). The accelerating voltage regulation range is 20–200 kV and the maximum magnification can exceed 1 000 000 times. The point resolution and line resolution can reach 0.24 nm and 0.102 nm, respectively. The glass samples for FE-TEM are ground into fine powder, and then dispersed in absolute ethanol. The density measurement of the samples is Archimedes method, using a high-precision balance with special accessories. UV-VIS-NIR transmission spectrum of glass was measured by a spectrophotometer (JASCO V-570). The region was 400–2 500 nm and the step length was 2 nm. FT-IR spectroscopy was measured by BRUKER VERTEX

70 and the region was 2–23 μm . The MIR fluorescence spectrum was measured by a fluorescence spectrometer (FSP920; Edinburgh Instruments Ltd., UK) equipped with a liquid-nitrogen cooled InSb detector (C4159-5671; HAMAMATSU, JPN). The measurement wavelength was 2 300–5 500 nm and the spectral step size was 2 nm, pumped by a 2 μm CW laser.

3 Results and Discussion

3.1 Amorphous and Thermal Properties

Tab. 1 is the EPMA analytical compositions of the undoped GAGSE glass sample after chemical purification treatment. Compared with original chemical composition, the deviation is so small that it can be accepted; thus, we regarded that the purified sample remains the original composition.

Fig. 2 shows the XRD patterns of all Pr^{3+} ions doped glass samples. Even the sample doped with high concentration Pr^{3+} ions as well as 0.4% has no obvious crystallization peaks. In order to make sure of amorphous properties and verify the dispersion of RE ions, the TEM test was further carried out, and the results are shown in Fig. 3 and 4. In both HR-TEM images of GAGSE-0.2 and GAGSE-0.3 samples, no crystal fringe is observed, indicating that the heavily Pr^{3+} ions doped samples still maintain the amorphous state. From the EDS results, it is obvious that all elements are dispersed evenly in the matrix. Comparing the EDS results of GAGSE-0.2 and GAGSE-0.3, it can be seen that the Pr^{3+} gathers a little more densely in some local area when the Pr^{3+}

Tab. 1 The EPMA result of GAGSE-undoped glass after chemical purification treatment

Composition	Theoretical chemical composition/%	Measured composition/%	Error/%
Ge	12	11.23	1.32
As	20.8	20.23	0.83
Ga	4	3.60	2.68
Se	63.2	64.19	0.59
O	0	0.59 [*]	22.3
Si	0	0.16 [*]	26.5
Al	0	0 [*]	100

*The results of these three elements are inauthentic due to the value of error.

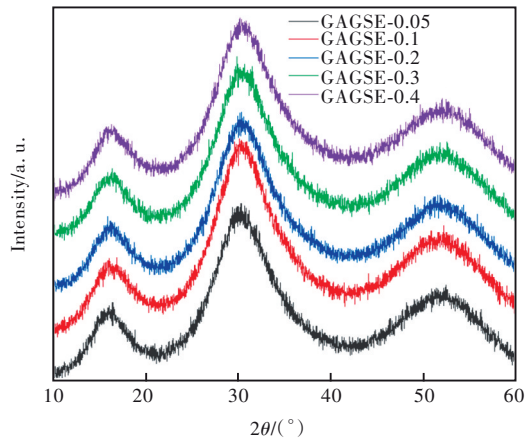


Fig. 2 XRD patterns of 0.05%–0.4% Pr³⁺ ions doped Ge₁₂-As_{20.8}Ga₄Se_{63.2} bulk samples

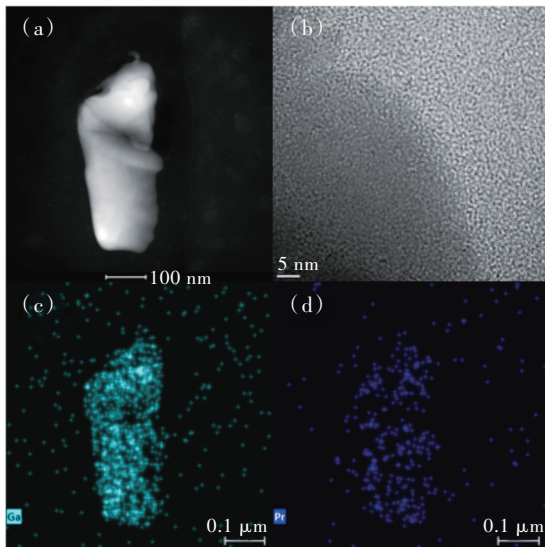


Fig. 3 (a) TEM image of GAGSE-0.2. (b) HR-TEM image of GAGSE-0.2. (c)–(d) The dispersion of Ga, Pr elements in GAGSE-0.2.

ions' concentration increases.

Tab. 2 shows the characteristic temperatures of all Pr³⁺ doped glasses. For the GAGSE-undoped glass, the transition temperature (T_g) is 212 °C and the initial crystallization temperature (T_x) is 336 °C.

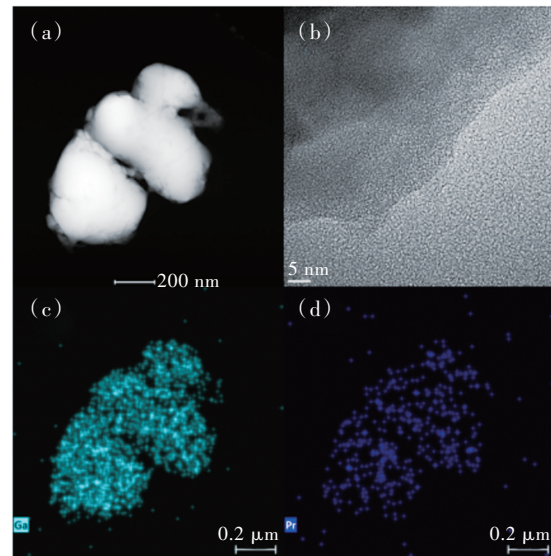


Fig. 4 (a) TEM image of GAGSE-0.3. (b) HR-TEM image of GAGSE-0.3. (c)–(d) The dispersion of Ga, Pr elements in GAGSE-0.3.

Tab. 2 Characteristic temperatures of 0–0.4% Pr³⁺ ions doped Ge₁₂As_{20.8}Ga₄Se_{63.2} glasses

Sample	$T_g/^\circ\text{C}$	$T_x/^\circ\text{C}$	$\Delta T/^\circ\text{C}$
GAGSE-undoped	212	336	124
GAGSE-0.1	215	334	119
GAGSE-0.2	217	331	114
GAGSE-0.3	220	330	110
GAGSE-0.4	223	328	105

Parameter of ΔT is used to evaluate the crystallization resistance and vitrification ability of glass. The larger value of ΔT indicates higher glass forming ability and better thermal stability. The ΔT of GAGSE-undoped glass is about 124 °C, indicating that the Ge₁₂As_{20.8}Ga₄Se_{63.2} glass has great thermal stability as a RE matrix glass. This is beneficial for fiber drawing in the next experiment. When the Pr³⁺ ions are introduced into the matrix glass, the T_g is slightly increased whereas T_x is decreased. It is obvious that

the GAGSE-0.4 glass still has a large ΔT value as well as 105 °C. Even the content of Pr^{3+} reaches 0.4%, the thermal stability of the glass is still suitable for fiber drawing.

3.2 Absorption and Emission Properties

Fig. 5 shows the absorption spectra of 0.05%–0.4% Pr^{3+} ions doped glasses in the wavelength range of 800–2500 nm and the electronic energy levels of Pr^{3+} ion. The four obvious absorption bands centered at 1064, 1494, 1608, 2036 nm correspond to the transitions from ground energy level of $^3\text{H}_4$ to the excited levels of $^1\text{G}_4$, $^3\text{F}_4$, $^3\text{F}_3$ and $^3\text{F}_2$, respectively. The cut-off edge is about 850 nm at the short wavelength because of the band gap absorption. The absorption coefficient of Pr^{3+} ions in $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ glass increases approximately linearly from 0.05% to 0.4% (see Fig. 6), which

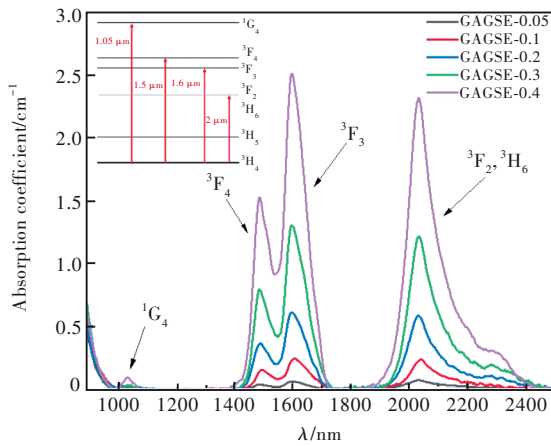


Fig. 5 Absorption spectra of 0.05%–0.4% Pr^{3+} ions doped $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ glasses in the range of 850–2500 nm

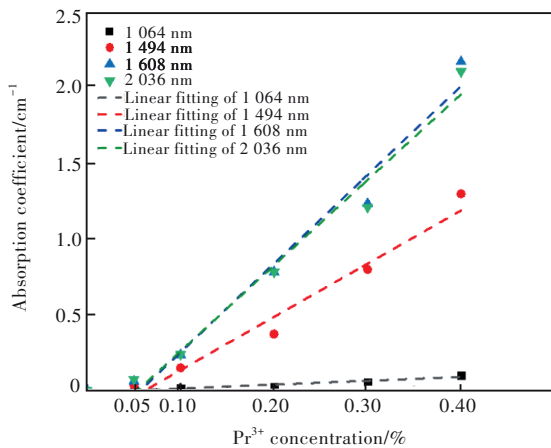


Fig. 6 Linear fitting of absorption coefficient of 0.05%–0.4% Pr^{3+} ions doped $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ glasses at each characteristic absorption peak

confirms the Pr^{3+} ions were dissolved almost homogeneously in the glass.

Based on the absorption spectrum and basic physical parameters, the Judd-Ofelt analysis is used to evaluate the luminescence characteristics^[19–20]. The oscillator strengths, Judd-Ofelt intensity parameters Ω_t ($t=2, 4, 6$) were calculated and summarized in Tab. 3. The values of other reported system glasses are also summarized for comparison. The Ω_t ($t=2, 4, 6$) are related to the local environment nearby the RE ions. Ω_2 is mainly affected by polarized and asymmetric environment around RE ions, Ω_4 is mainly affected by the acidity and alkaline, and Ω_6 is related to the rigidity of glass^[21–23]. In addition, the value of Ω_4/Ω_6 is used to evaluate the spontaneous emission characteristics of RE ions. The two ground state electronic absorption peaks of $^3\text{H}_4 \rightarrow ^3\text{F}_4$, $^3\text{F}_3$ overlap each other. For J-O analysis, this wide absorption peak can be divided into two characteristic peaks. Taking sample GAGSE-0.4 as an example, the result of division is shown in Fig. 7. The fitting mode is Gaussian-Lorentz and the R^2 is 0.997455. It can be found that the Ω_t suddenly changed when the concentration of Pr^{3+} ions increased to 0.3%. The sudden increase of Ω_2 implies the asymmetric environment was changed around RE ions. In general, the low symmetry of polyhedral structure and high covalency of constituent bonds will lead to a larger oscillator strength. Compared with the other hosts, Pr^{3+} ions doped GAGSE glass has the larger oscillator strength.

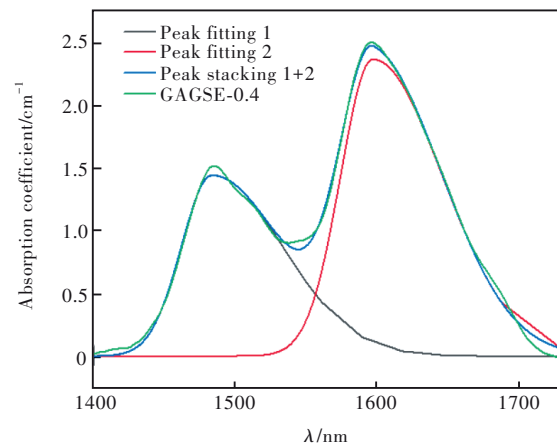


Fig. 7 The division fitting of $^3\text{H}_4 \rightarrow ^3\text{F}_4$, $^3\text{F}_3$ absorption peak in GAGSE-0.4 glass.

Tab. 3 The J-O intensity parameters and oscillator strength of 0.05%–0.4% Pr³⁺ doped Ge₁₂As_{20.8}Ga₄Se_{63.2} glasses and other hosts^[24–28]

Glass	J-O parameters/(10 ⁻²⁰ cm ²)				Oscillator strength/10 ⁻⁶				Reference
	Ω_2	Ω_4	Ω_6	Ω_4/Ω_6	³ H ₄ →				
					¹ G ₄	³ F ₄	³ F ₃	³ F ₂	
GAGSE-0.05	2.37	16.51	8.59	1.92	1.096	13.194	28.857	17.356	This work
GAGSE-0.1	2.59	14.63	6.05	2.42	1.009	9.712	22.948	15.159	
GAGSE-0.2	2.5	14.15	5.86	2.41	1.074	9.412	22.22	14.663	
GAGSE-0.3	2.98	11.37	4.77	2.38	0.812	7.811	18.411	12.965	
GAGSE-0.4	3.3	10.17	4.73	2.15	0.764	7.497	17.136	12.8	[24]
Phosphate	4.26	4.33	6.27	0.69	0.250	10.540	10.640	5.000	
Fluorotellurite	3.57	6.60	5.18	1.27	0.409	4.543	9.668	6.357	
ZBLAN	0.94	6.54	3.84	1.70	3.14	4.05	2.09	0.99	
BKZ	4.92	14.08	8.48	1.66	—	5.49	12.08	7.89	[27]
GAGS	10.4	4.78	10.4	0.46	—	—	—	—	[28]

Fig. 8 shows the infrared transmission spectra of blank and 0.05%–0.4% Pr³⁺ ions doped GAGSE glasses. The infrared cut-off wavelengths are about 18 μm. Because the raw materials have been purified, there are no obvious peaks of impurity absorptions before 9 μm. The strongest absorption peak is at around 4.6 μm which is caused by the absorption of ³H₅ energy level of Pr³⁺ ions. However, the absorption of Se—H vibration (3.5 μm and 4.12 μm) with the introduction of RE ions cannot be completely excluded. Besides, the weak absorption peak at 6.3 μm is caused by the impurity absorption of H₂O.

In addition, comparing the Pr³⁺ ions doped and undoped glasses' spectra, it can be found that the Se

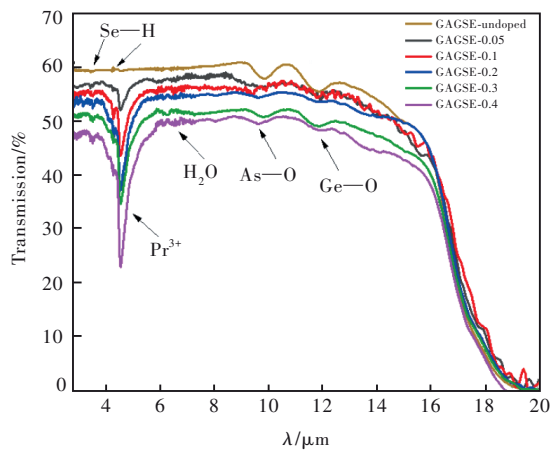


Fig.8 Infrared transmission spectra of 0.05%–0.4% Pr³⁺ ions doped Ge₁₂As_{20.8}Ga₄Se_{63.2} glasses compared with an undoped one in the range of 2–20 μm (The transmission results in this figure are adjusted by multiplying the coefficient for display)

—H vibration at 3.5 μm and 4.12 μm will appear due to the introduction of Pr³⁺ ions. This is probably because the raw material of Praseodymium powder was generally prepared by a hydrogenation method, which will result in a lot of H elements absorbed in the powder can hard to be totally removed through the subsequent processing process. Moreover, the strong absorption peaks at 9.7 μm and 11.8 μm should be caused by the vibration of As—O and Ge—O, respectively^[17]. In fact, it can be seen from the experimental results that the present purified process has an effect on the removal of Se—H, —OH and H₂O, but has no obvious work for the elimination of As—O and Ge—O.

Fig. 9 shows the MIR emission spectra of 0.05%–0.4% Pr³⁺ ions doped Ge₁₂As_{20.8}Ga₄Se_{63.2}

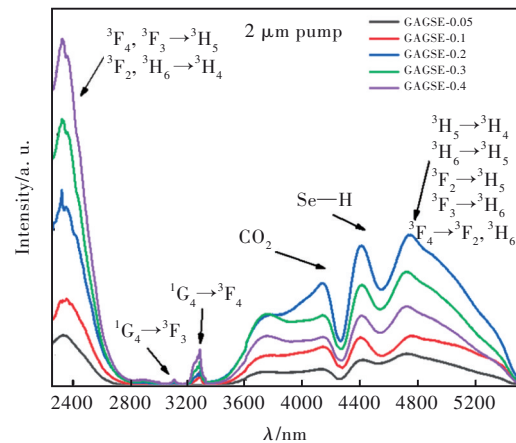


Fig.9 Mid-infrared emission spectra of 0.05%–0.4% Pr³⁺ ions doped Ge₁₂As_{20.8}Ga₄Se_{63.2} glasses pumped by a 2.0 μm laser

glasses ranging from 2 300 nm to 5 500 nm pumped by a 2.0 μm laser. The Pr^{3+} ions have a very wide emission band from 3.4 nm to 5.5 μm . The emission peaks, which centered at 2.33, 3.1, 3.3 μm , correspond to the transitions of ${}^3\text{F}_2/{}^3\text{H}_6 \rightarrow {}^3\text{H}_4$, ${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$, and ${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$, respectively. The electrons that have already been excited to the ${}^3\text{F}_2$ energy level will continue to be excited to the higher ${}^1\text{G}_4$ energy level when pumped by a 2.0 μm laser, as shown in Fig.

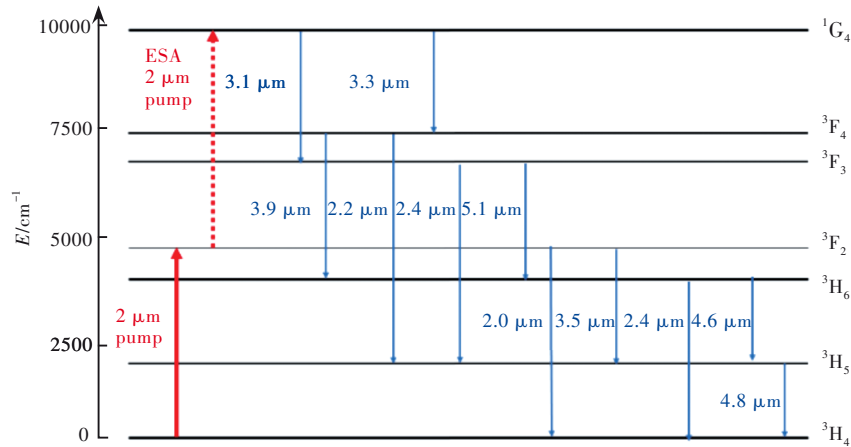


Fig.10 The schematic energy level diagram of Pr^{3+} ions

From the fluorescence spectral variation, it can be seen that when the Pr^{3+} ions' concentration is lower than 0.2%, the fluorescence intensity increases with the increasing of the Pr^{3+} ions' concentration and reaches the maximum at 0.2%. It can also be found that further increasing of Pr^{3+} ions' concentration will enhance the intensities of emissions at 2.33 μm and 3.1 μm but the intensity of wide emission band (3.4–5.5 μm) is decreased. This phenomenon indicates that when the concentration of Pr^{3+} ions is less than or equal 0.2%, most of the Pr^{3+} ions can be evenly distributed in the glass matrix by mainly forming Pr—Se—Ga bonds. However, when the concentration of Pr^{3+} ions is further increased, the finite non-bridging Se can no longer provide enough negative coordination for the Pr^{3+} ions. Therefore, the Pr^{3+} ions gradually begin to be clustered (see Fig. 3(d)) to share ligands by forming of Pr—Se—Pr instead of Pr—Se—Ga ones. In this way, the part of Pr^{3+} ions' clusters contribute little to MIR fluorescence, which ultimately leads to fluorescence quenching^[9,30]. Therefore, according to the J-O

10. Actually, the emission peak at 2.33 μm is also contributed by the transition of ${}^3\text{F}_4/{}^3\text{F}_3 \rightarrow {}^3\text{H}_5$ ^[29]. The broadband luminescence is contributed by the superposition of emission peaks (*i. e.*, ${}^3\text{F}_4 \rightarrow {}^3\text{F}_2/{}^3\text{H}_6$, ${}^3\text{F}_3 \rightarrow {}^3\text{H}_6$, ${}^3\text{F}_2 \rightarrow {}^3\text{H}_5$, ${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$ and ${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$). In addition, the dip at 4.24 μm is ascribed to the absorption of CO_2 in the air of the spectrometer and the dip at 4.55 μm is caused by the Se—H bond impurity of the sample itself^[8-9,30].

parameters and fluorescence spectra, it is concluded that the optimal doping concentration of Pr^{3+} ions in this glass matrix is 0.2%. When the concentration is higher than 0.2%, the clusters will bring about fluorescence quenching by depopulating the excited states.

The stimulated emission cross-section of Pr^{3+} in this glass matrix is calculated basing on the Füchtbauer-Ladenburg equation^[31]:

$$\sigma_{\text{emi}}(\lambda) = \frac{\lambda_p^4 \lambda I(\lambda)}{8\pi c n^2 \tau_{\text{rad}} \int \lambda I(\lambda) d\lambda}, \quad (1)$$

in which the effective line-width, λ_{eff} , of fluorescence band is defined by:

$$\Delta\lambda_{\text{eff}} = \frac{\int I(\lambda) d\lambda}{I_{\text{max}}}, \quad (2)$$

the $\int I(\lambda) d\lambda$ is the effective peak area and I_{max} is the maximum intensity of each emission peak. λ_p is the center wavelength of emission peak. c is speed of light, n is refractive index, τ_{rad} is the radiation lifetime of transition. The calculated results are listed in Tab. 4. The GAGSE-0.2 glass has a large σ_{emi} ,

long upper state lifetime, and good branching ratios for transition ${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$ and ${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$. This is mainly because the chalcogenide glass has a larger refrac-

tive index and higher degree of covalent bonding. These characteristics of GAGSE-0.2 glass make it possible for MIR fiber amplifiers and lasers.

Tab. 4 The radiative parameters of 0.05%–0.4% Pr³⁺ doped Ge₁₂As_{20.8}Ga₄Se_{63.2} glasses calculated from J-O parameters

Class	Transition	$\lambda_{\text{emi}}/\text{nm}$	$\lambda_{\text{eff}}/\text{nm}$	$\tau_{\text{rad}}/\text{ms}$	$\beta/\%$	$A_{\text{rad}}/\text{s}^{-1}$	$\sigma_{\text{emi}}/(10^{-20}\text{cm}^2)$
GAGSE-0.05	${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$	3 434	131.22	5.92	1.86	133.94	2.491
	${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$	3 000	82.63	—	0.48	34.88	0.677
	${}^3\text{F}_3 \rightarrow {}^3\text{H}_5$	2 370	—	0.80	15.41	1 243.76	—
	${}^3\text{F}_4 \rightarrow {}^3\text{H}_5$	2 150	—	0.62	35.71	1 592.57	—
	${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$	4 814	752.79	8.59	100	116.36	1.582
	${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$	4 600	734.94	8.16	43.26	122.43	1.401
GAGSE-0.1	${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$	3 434	137.14	8.18	1.68	96.49	1.717
	${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$	3 000	78.80	—	0.44	25.70	0.523
	${}^3\text{F}_3 \rightarrow {}^3\text{H}_5$	2 370	—	0.89	17.06	1 116.32	—
	${}^3\text{F}_4 \rightarrow {}^3\text{H}_5$	2 150	—	0.77	37.87	1 290.25	—
	${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$	4 814	743.23	10.90	100	91.66	1.305
	${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$	4 600	730.65	10.32	44.16	96.82	1.114
GAGSE-0.2	${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$	3 434	151.73	8.45	1.68	93.43	1.492
	${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$	3 000	102.36	—	0.45	24.88	0.333
	${}^3\text{F}_3 \rightarrow {}^3\text{H}_5$	2 370	—	0.92	17.04	1 079.16	—
	${}^3\text{F}_4 \rightarrow {}^3\text{H}_5$	2 150	—	0.80	37.86	1 248.55	—
	${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$	4 814	741.76	11.25	100	88.84	1.193
	${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$	4 600	722.42	10.65	44.19	93.83	1.066
GAGSE-0.3	${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$	3 434	171.31	10.04	1.73	79.19	1.120
	${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$	3 000	100.24	—	0.44	20.43	0.279
	${}^3\text{F}_3 \rightarrow {}^3\text{H}_5$	2 370	—	1.03	18.45	968.83	—
	${}^3\text{F}_4 \rightarrow {}^3\text{H}_5$	2 150	—	0.98	37.50	1 013.12	—
	${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$	4 814	728.76	13.51	100	77.94	1.011
	${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$	4 600	752.17	12.82	44.81	93.83	0.850
GAGSE-0.4	${}^1\text{G}_4 \rightarrow {}^3\text{F}_4$	3 434	163.56	10.04	1.83	79.66	1.180
	${}^1\text{G}_4 \rightarrow {}^3\text{F}_3$	3 000	103.23	—	0.46	19.98	0.265
	${}^3\text{F}_3 \rightarrow {}^3\text{H}_5$	2 370	—	1.07	18.77	933.02	—
	${}^3\text{F}_4 \rightarrow {}^3\text{H}_5$	2 150	—	1.05	36.39	945.62	—
	${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$	4 814	719.76	14.06	100	71.10	0.984
	${}^3\text{H}_6 \rightarrow {}^3\text{H}_5$	4 600	749.27	13.39	44.85	74.63	0.817

3.3 Fiber Fabrication and Loss

The GAGSE-0.2 glass sample was selected and drawn into a single-mode, double-cladding fiber. Fig. 11 shows the sectional micrographs at different magnifications and the outermost layer is FEP plastic. The diameter of core, inner cladding and outer cladding are 15, 140, 170 μm , respectively. The refractive indexes ($n_{\text{core}}=2.613$, $n_{\text{inner-cladding}}=2.611$, $n_{\text{outer-cladding}}=2.587$) were measured by an IR VASE Mark II Spectroscopic Ellipsometer at 2 μm wavelength. The numerical apertures (N. A.) between core and

inner cladding is about 0.10, between inner and outer cladding is about 0.35, which were calculated at 2 μm wavelength. There is no obvious crystallization can be found on the surface of fiber, indicating this glass has good ability for fiber fabrication.

The fiber loss of GAGSE-0.2 double-cladding fiber is shown in Fig. 12. The test method is truncation method. The fiber has a lowest loss of 2.95 dB/m at 6.58 μm . The absorption band at 4.6 μm is due to the Pr³⁺ (${}^3\text{H}_5 \rightarrow {}^3\text{H}_4$) and Se—H bonds. The absorption bands at 3.5, 6.3, 7.9 μm should be

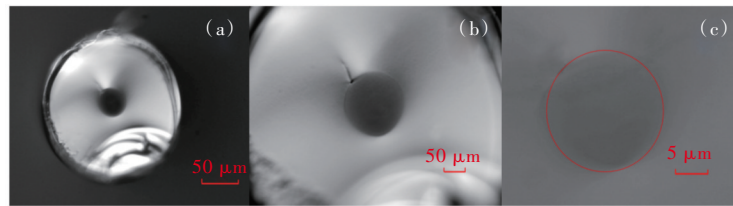


Fig.11 The sectional micrographs of 0.2% Pr^{3+} -doped $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ double-cladding fiber. (a)50× magnified image of whole fiber. (b)100× magnified image of outer cladding and inner cladding. (c)500× magnified image of core and inner cladding.

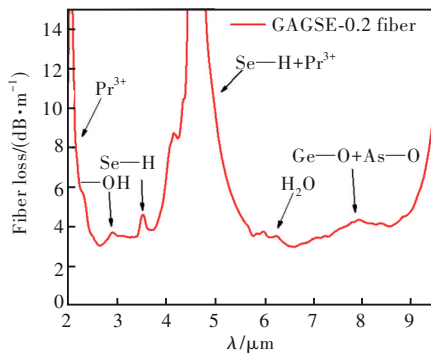


Fig.12 The optical loss of 0.2% Pr^{3+} -doped $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ double-cladding fiber

—OH, Se—H, H_2O and oxide impurities (Ge—O, As—O), respectively.

4 Conclusion

The new glass composition in the Ge-As-Ga-Se chalcogenide system, *i. e.* $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ shows high Pr^{3+} ions' solubility (0.4% in weight at maximum)

and good thermal characteristic ($T_x - T_g \geq 105$ °C). The Pr^{3+} ions have a very wide emission band from 3.4 μm to 5.5 μm and the 0.2% (in weight) Pr^{3+} ions doped glass shows the best mid-infrared emission characteristic. Through physical and chemical purification, the impurity content can be effectively reduced to an acceptable level. The 0.2% (in weight) Pr^{3+} ions doped $\text{Ge}_{12}\text{As}_{20.8}\text{Ga}_4\text{Se}_{63.2}$ glass was successfully drawn into a step-index, double-cladding fiber by a multistage rod-in-tube method, and the lowest loss is 2.95 dB/m at 6.58 μm . To sum up, it should be a good potential gain medium to realize mid-infrared fiber lasers working at 3.5–5.5 μm .

Response Letter is available for this paper at: <http://cjl.lightpublishing.cn/thesisDetails#10.37188/CJL.20220088>.

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